

# Vertex operators, Kronecker products, and Hilbert series

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# Summary of yesterday's lecture

## 1 Hopf:

- $\Delta f = f(X + Y)$
- $\langle f \cdot g, h \rangle = \langle f \otimes g, \Delta h \rangle$
- $\text{ch}(\chi^\lambda) = s_\lambda$
- product = induction, coproduct = restriction

## 2 Vertex:

$$\sum_{n \in \mathbb{Z}} s_{(n, \nu)} = \Gamma_1 s_\nu = \sigma_1 D_{\lambda_{-1}} s_\nu$$

# Summary of today's lecture

- 1 Kronecker
- 2 Kronecker + Hopf
- 3 Reduced notation = Kronecker + Hopf + Vertex
- 4 Application to Hilbert series of some invariant algebras

# The internal product of $Sym$

- Can be defined without reference to characters
- Remember Cauchy's identity

$$\sigma_1(XY) = \sum_{\lambda} s_{\lambda}(X)s_{\lambda}(Y)$$

- $\delta : f \mapsto f(XY)$  is a coproduct
- Obviously,  $\delta p_{\mu} = p_{\mu} \otimes p_{\mu}$
- The dual product is  $p_{\mu} * p_{\nu} = z_{\mu} \delta_{\mu\nu} p_{\mu}$
- It corresponds under  $ch$  to the pointwise product of class functions

- 

$$s_{\mu} * s_{\nu} = \sum_{\lambda} g_{\mu\nu}^{\lambda} s_{\lambda} = ch(\chi^{\mu} \chi^{\nu})$$

# The splitting (or Mackey) formula I

- There is a compatibility between  $*$ ,  $\cdot$  and  $\Delta$
- It reflects a general formula in group theory
- Again, it has a direct and elementary proof
- This is

$$(fg) * h = \mu[(f \otimes g) * \Delta h]$$

where  $\mu(u \otimes v) = uv$  and

$$(a \otimes b) * (a' \otimes b') = (a * a') \otimes (b * b')$$

- Generalization  $(f_1 f_2 \cdots f_r) * h$

## The splitting (or Mackey) formula II

**Proof** (Hopf style):

$$\begin{aligned} \langle (fg) * h, u \rangle &= \langle (fg)(X)h(Y), u(XY) \rangle \\ &= \langle f(X')g(X'')h(Y), u(X'Y + X''Y) \rangle \\ &= \sum_{(u)} \langle f(X')g(X'')h(Y), u_{(1)}(X'Y)u_{(2)}(X''Y) \rangle \end{aligned}$$

(the right part is a  $Y$  product that we can dualize)

$$\begin{aligned} &= \sum_{(u)} \langle f(X')g(X'')h(Y' + Y''), u_{(1)}(X'Y')u_{(2)}(X''Y'') \rangle \\ &= \sum_{(h)} \langle f(X')g(X'')h_{(1)}(Y')h_{(2)}(Y''), u(X'Y' + X''Y'') \rangle \end{aligned}$$

## The splitting (or Mackey) formula III

$$= \sum_{(h)} \langle (f * h_{(1)})(X') (g * h_{(2)})(X''), u(X' + X'') \rangle$$

(now  $X' Y' \rightarrow X'$  and  $X'' Y'' \rightarrow X''$ )

$$= \langle \mu[(f \otimes g) * \Delta h, u] \rangle .$$

**Example:**

$$h_\mu * h_\nu = \sum_{M \in \mathcal{M}(\mu, \nu)} h_M$$

# The reduced notation I

- Murnaghan, Littlewood:

$$\langle \mu \rangle * \langle \nu \rangle = \sum_{\lambda} \bar{g}_{\mu\nu}^{\lambda} \langle \lambda \rangle$$

means

$$s_{\mu[n]} * s_{\nu[n]} = \sum_{\lambda} \bar{g}_{\mu\nu}^{\lambda} s_{\lambda[n]}$$

- But what is  $\langle \lambda \rangle$ , precisely ?
- Answer: image of  $s_{\lambda}$  by the vertex operator

$$\langle \lambda \rangle = \Gamma_1 s_{\lambda} = \sum_{m \in \mathbb{Z}} s_{(m, \lambda)}$$

That is, a generating series ...



## The reduced notation II

- This follows from

$$(\sigma_1 f) * (\sigma_1 g) = \sigma_1 \sum_{\mu, \nu} (D_{v_\mu} f)(D_{v_\nu} g)(u_\mu * u_\nu)$$

where  $(u, v)$  is any pair of adjoint bases of  $Sym$

**Proof:**

$$\begin{aligned} (\sigma_1 f) * (\sigma_1 g) &= \mu[(\sigma_1 \otimes f) * \Delta \sigma_1 \Delta g] \\ &= \mu[(\sigma_1 \otimes f) * \left( \sum_{\gamma} D_{v_\gamma} g \otimes u_\gamma \right) (\sigma_1 \otimes \sigma_1)] \\ &= \mu[(\sigma_1 \otimes f) * \left( \sum_{\gamma} \sigma_1 D_{v_\gamma} g \right) \otimes \sigma_1 u_\gamma] \\ &= \sum_{\gamma} \mu[(\sigma_1 D_{v_\gamma} g) \otimes (f * \sigma_1 u_\gamma)] \end{aligned}$$

## The reduced notation III

$$\begin{aligned}
 &= \sum_{\gamma} (\sigma_1 D_{v_{\gamma}} g) \mu [(\sigma_1 \otimes u_{\gamma}) * \sum_{\delta} D_{v_{\delta}} f \otimes u_{\delta}] \\
 &= \sum_{\gamma, \delta} (\sigma_1 D_{v_{\gamma}} g) (D_{v_{\delta}} f) (u_{\gamma} * u_{\delta}).
 \end{aligned}$$

Applying this to  $u = v = s$  and  $f = s_{\mu}(X - 1)$ ,  $g = s_{\nu}(X - 1)$ , we get Littlewood's formula, which reads now

$$\Gamma_1 s_{\mu} * \Gamma_1 s_{\nu} = \sum_{\lambda} \bar{g}_{\mu\nu}^{\lambda} \Gamma_1 s_{\lambda}$$

or, more explicitly

$$\Gamma_1 s_{\mu} * \Gamma_1 s_{\nu} = \Gamma_1 \sum_{\alpha\beta\gamma} (D_{s_{\gamma}} D_{s_{\alpha}} s_{\mu}) (D_{s_{\gamma}} D_{s_{\beta}} s_{\nu}) (s_{\alpha} * s_{\beta})$$

## The reduced notation IV

**Example:** Two-row shapes,  $s_{(n-k,k)} * s_{(n-l,l)}$

$$\Gamma_1 s_k * \Gamma_1 s_l = \Gamma_1 \sum_{p=0}^{\min(k,l)} \sum_{q=0}^p s_{k-p} s_{l-p} s_{p-q}$$

The triple product of one-part Schur functions is easily evaluated. With  $k = 2, l = 3$ , we get

$$\Gamma_1 (s_2 s_3 s_0 + s_1 s_2 (s_1 + s_0) + s_0 s_1 (s_2 + s_1 + s_0))$$

$$= \Gamma_1 (s_{32} + s_{41} + s_5 + s_{211} + s_{22} + 2s_{31} + s_4 + 2s_{21} + 2s_3 + s_{11} + s_2 + s_1)$$

so that

$$s_{82} * s_{73} = s_{532} + s_{541} + s_{55} + s_{6211} + s_{622} + s_{631} + s_{64} + 2s_{721} + 2s_{73} + s_{811} + s_{82} + s_{91}$$

Schur functions and  $GL(n, \mathbb{C})$ 

- I. Schur (1901) The irreducible polynomial representations of  $GL(n, \mathbb{C})$  are parametrized by partitions in at most  $n$  parts
- if  $V = \mathbb{C}^n$  the representations of degree  $k$  are those occurring in  $V^{\otimes k}$

$$V^{\otimes k} = \bigoplus_{\lambda \vdash k, l(\lambda) \leq n} V_{\lambda}^{\oplus f_{\lambda}}$$

( $f_{\lambda}$  = nb of standard tableaux of shape  $\lambda$ )

- Character formula:

$$\text{tr } \rho_{\lambda}(g) = s_{\lambda}(g)$$

(as a symmetric function of the eigenvalues of  $g$ )

- Examples: symmetric tensors =  $h_k$ , alternating tensors =  $e_k$ , determinant =  $e_n$
- Proof: Schur-Weyl duality

# Schur functions and $SL(n, \mathbb{C})$

- $V_\lambda$  remains irreducible, but now  $e_n = 1$
- So  $V_{\lambda+(1^n)} \simeq V_\lambda$
- In particular,  $V_{(m^n)}$  is the trivial representation
- Invariants of  $SL(n)$  come from rectangular shapes

# Invariants of multilinear forms I

- Irreducible representations of a product group  
 $G = \prod_{i=1}^k GL(n_i)$ : the characters are  $\prod_{i=1}^k s_{\lambda^{(i)}}(X_i)$
- We are interested in the relative invariants of  $G$  in  $S^d(V_1 \otimes \cdots \otimes V_k)$ , where  $V_i = \mathbb{C}^{n_i}$ , i.e., homogeneous polynomials  $F$  in the coordinates such that

$$g \cdot F = (\det g_1)^{l_1} (\det g_2)^{l_2} \cdots (\det g_k)^{l_k} F$$

for any  $g = (g_1, \dots, g_k) \in G$

- A *covariant* of degree  $d = (d_0, d_1, \dots, d_k)$  is a relative invariant of  $G$  in the representation space

$$S^{d_0}(V_1 \otimes \cdots \otimes V_k) \otimes S^{d_1}(V_1^*) \otimes \cdots \otimes S^{d_k}(V_k^*)$$

## Invariants of multilinear forms II

If we write  $A \in V_1^* \otimes \cdots \otimes V_k^*$  as

$$A(\mathbf{x}_1, \dots, \mathbf{x}_k) = A_{i_1 i_2 \dots i_k} x_1^{i_1} x_2^{i_2} \cdots x_k^{i_k}$$

The coordinate functions  $x_j^i, j = 1, \dots, n_j$ , form a basis of  $V_j^*$  and the components  $A_{i_1 \dots i_k}$  are regarded as a basis of  $(V_1^* \otimes \cdots \otimes V_k^*)^* = \mathbf{V}$ . An invariant  $F$  is a homogeneous polynomial in the coefficients of the “groundform”  $A$ , such that  $F = 0$  defines a  $G$ -invariant hypersurface of  $\mathbb{P}(\mathcal{V})$ . Similarly, a covariant is a multi-homogeneous polynomial in the original vector variables  $\mathbf{x}_j$ , whose coefficients are homogenous polynomials in the  $A_{i_1 \dots i_k}$ , of which the simultaneous vanishing defines a  $G$ -invariant subvariety of  $\mathbb{P}(\mathbf{V})$ .

# Invariants of multilinear forms III

A covariant is a  $G$ -equivariant map from  $S^{d_0}(V_1 \otimes \cdots \otimes V_k)$  to the irreducible representation  $S^{d_1}(V_1) \otimes \cdots \otimes S^{d_k}(V_k)$ . In general, a concomitant of degree  $d_0$  and of type  $\lambda = (\lambda^{(1)}, \dots, \lambda^{(k)})$ , where the  $\lambda^{(i)}$  are partitions, is an equivariant map from the same space to the irreducible representation  $S_{\lambda^{(1)}}(V_1) \otimes \cdots \otimes S_{\lambda^{(k)}}(V_k)$  of  $G$ .



## Invariants of multilinear forms IV

Characters of the irreducible polynomial representations of the product group  $G$

$$S_\lambda = s_{\lambda^{(1)}}(X_1) \cdots s_{\lambda^{(k)}}(X_k)$$

$\lambda^{(i)}$  are partitions,  $X_i = \{x_{i1}, \dots, x_{in_i}\}$  is a set of  $n_i$  variables.  
The character of the one dimensional representation

$$\det^l(g) = (\det g_1)^{l_1} (\det g_2)^{l_2} \cdots (\det g_k)^{l_k}$$

is the product of rectangular Schur functions

$$s_{(l_1^{n_1})}(X_1) s_{(l_2^{n_2})}(X_2) \cdots s_{(l_k^{n_k})}(X_k)$$

The character of  $G$  in  $S^d(\mathbf{V})$  is  $h_d(X_1 X_2 \cdots X_k)$ .

# Invariants of multilinear forms V

Hence, the dimension of the space of invariants of degree  $d$  and weight  $l$ , which is also the multiplicity of the one dimensional character  $\det^l$  in  $S^d(\mathbf{V})$ , is

$$\dim \text{Inv}(d, l) = \langle h_d(X_1 X_2 \cdots X_k), s_{(l_1^{n_1})}(X_1) s_{(l_2^{n_2})}(X_2) \cdots s_{(l_k^{n_k})}(X_k) \rangle_G$$

Replace the  $X_i$  by infinite sets of independent variables, and compute in  $\text{Sym}^{\otimes k}$  is dual to the internal product  $*$

$$\begin{aligned} \dim \text{Inv}(d, l) &= \langle \delta^k(h_d), s_{(l_1^{n_1})} \otimes \cdots \otimes s_{(l_k^{n_k})} \rangle_{\text{Sym}^{\otimes k}} \\ &= \langle h_d, s_{(l_1^{n_1})} * \cdots * s_{(l_k^{n_k})} \rangle_{\text{Sym}}. \end{aligned}$$

# Invariants of multilinear forms VI

The internal product of two homogenous symmetric functions being zero if these are not of the same degree, we see that  $\text{Inv}(d, l)$  can be nonzero only if the conditions

$$n_1 l_1 = n_2 l_2 = \cdots = n_k l_k = d$$

are satisfied. In particular, if all the  $n_i$  are equal, the  $l_i$  must also be all equal.

## Invariants of multilinear forms VII

Let  $c(d; l)$  be the dimension of the space of covariants of degree  $d = (d_0, d_1, \dots, d_k)$  and weight  $l = (l_1, \dots, l_k)$ .

$$\begin{aligned} c(d; l) &= \langle h_{d_0}(X_1 X_2 \cdots X_k), (s_{(l_1^{n_1})} h_{d_1})(X_1) \cdots (s_{(l_k^{n_k})} h_{d_k})(X_k) \rangle_G \\ &= \langle h_{d_0}, (s_{(l_1^{n_1})} h_{d_1}) * \cdots * (s_{(l_k^{n_k})} h_{d_k}) \rangle_{\text{Sym}}. \end{aligned}$$

For  $SL(2)$ ,  $s_{(l,l)} h_d = s_{(l+d,l)}$ , so that the covariants are in bijection with highest weight vectors.

## Multilinear binary forms (qubit systems) I

If all  $V_i = \mathbb{C}^2$  we need only two-part partitions  
For the size  $(2, 2, 2)$ , we have

$$\dim \text{Inv}(2l; l, l, l) = \langle h_{2l}, s_{ll}^{*3} \rangle = \langle s_{ll} * s_{ll}, s_{ll} \rangle = \begin{cases} 0 & l \text{ odd} \\ 1 & l \text{ even} \end{cases}$$

using first the property  $\langle f * g, h \rangle = \langle f, g * h \rangle$  and the formula for  $s_{ll} * s_{ll}$ . Hence,

$$\sum_{d \geq 0} \dim S^d(\mathbf{V}) G t^d = \frac{1}{1 - t^4}.$$

The algebra of invariants is in this case  $\mathbb{C}[\Delta]$ , where  $\Delta = \text{Det}(A)$  is the hyperdeterminant.

## Multilinear binary forms (qubit systems) II

The generating series for the covariants can be written in the form

$$C(t; \mathbf{u}; \mathbf{v}) = \sum_{d,l} c(d, l) t^{d_0} u_1^{d_1} u_2^{d_2} u_3^{d_3} v_1^{l_1} v_2^{l_2} v_3^{l_3}$$

$$= \langle \sigma_1[tu_1s_1 + t^2v_1s_{11}], \sigma_1[u_2s_1 + v_2s_{11}] * \sigma_1[u_3s_1 + v_3s_{11}] \rangle$$

since with two variables,

$$\sigma_1[vs_{11}] = \sum_{l \geq 0} v^l s_{ll}$$

and

$$\sigma_1[us_1 + vs_{11}] = \sum_{\ell(\lambda) \leq 2} u^{\lambda_1 - \lambda_2} v^{\lambda_2} s_\lambda(X).$$

## Multilinear binary forms (qubit systems) III

The last sum can be obtained by combining vertex operators and MacMahon's linear operator  $\Omega_{\geq}^u$ , which maps any monomial containing a negative power of  $u$  to 0.

$$\begin{aligned} \sum_{\lambda_1 \in \mathbb{Z}, \lambda_2 \geq 0} t^{|\lambda|} u^{\lambda_1 - \lambda_2} v^{\lambda_2} s_{\lambda}(X) &= \Gamma_{tu} \Gamma_{tv/u}(1) \\ &= \left(1 - \frac{v}{u^2}\right) \sigma_t \left[ \left(u + \frac{v}{u}\right) X \right]. \end{aligned}$$

Hence, if  $\Omega_{\geq}^{\mathbf{u}}$  denotes the MacMahon operator annihilating any monomial containing a negative power of any of the  $u_i$ ,

$$C(t; \mathbf{u}; \mathbf{v}) = \Omega_{\geq}^{\mathbf{u}} \prod_{i=1}^3 \left(1 - \frac{v_i}{u_i^2}\right) \sigma_t \left[ \prod_{i=1}^3 \left(u_i + \frac{v_i}{u_i}\right) \right].$$

## Multilinear binary forms (qubit systems) IV

Here,  $\Omega$  is easily computed with the help of a computer algebra system by decomposing the right-hand side into partial fractions, and throwing away the terms leading to negative powers of the  $u_i$  in the Laurent expansion.

Setting the  $v_i$  equal to 1, one finds

$$\frac{1 - t^6 u_1^2 u_2^2 u_3^2}{(1 - t u_1 u_2 u_3)(1 - t^2 u_1^2)(1 - t^2 u_2^2)(1 - t^2 u_3^2)(1 - t^3 u_1 u_2 u_3)(1 - t^4)}$$

The structure of the generating series is the same for  $k$  qubits:

$$C(t; \mathbf{u}; \mathbf{v}) = \Omega_{\geq}^{\mathbf{u}} \prod_{i=1}^k \left( 1 - \frac{v_i}{u_i^2} \right) \sigma_t \left[ \prod_{i=1}^k \left( u_i + \frac{v_i}{u_i} \right) \right].$$



# Multilinear binary forms (qubit systems) V

For  $k = 4$ , the result is huge, and can be obtained only with more subtle algorithms (e.g., Xin's), but setting  $u_i = 0$  after each  $\Omega_{\geq}^{u_i}$  gives easily the Hilbert series of invariants.

For  $k = 5$ , this still works for the invariants.

Similar (but harder) calculations would give the Hilbert series of unitary or special unitary invariants.