

Crystal energy via charge in types A and C

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based on arXiv:1107.4169 (Math. Zeitschrift)
and work joint with **Naito**, **Sagaki**, **Shimozono** (in progress)

Outline

Crystals

Energy function

Charge

Arbitrary type

Outline

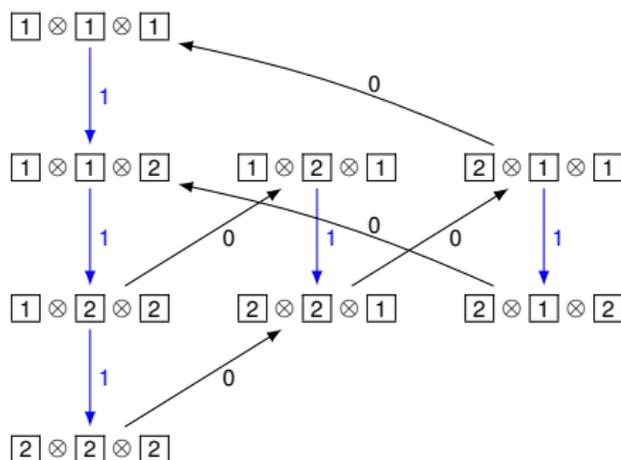
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Crystal graph



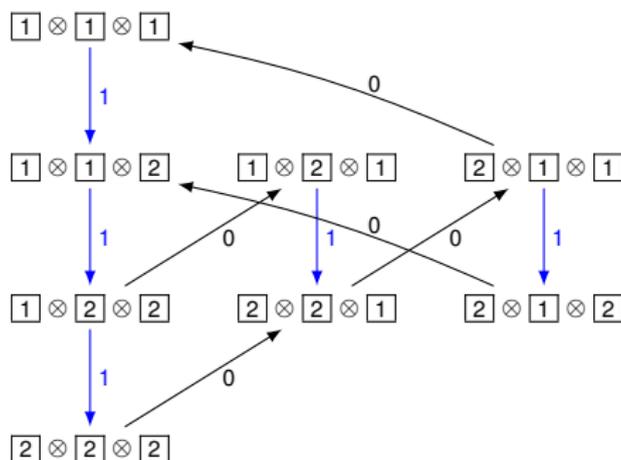
A $U_q(\mathfrak{g})$ -crystal is a nonempty set B with maps

$$\text{wt}: B \rightarrow P$$

$$e_i, f_i: B \rightarrow B \cup \{\emptyset\} \quad \text{for all } i \in I$$

Write $b \xrightarrow{i} b'$ for $b' = f_i(b)$

Crystal graph



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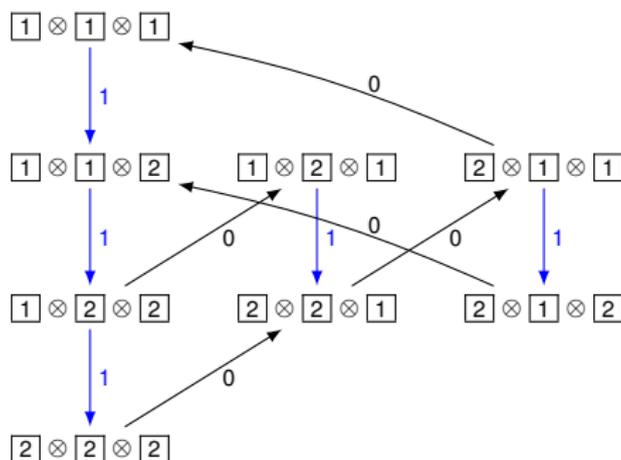
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Kashiwara–Nakashima tableaux

embed $B(1^N) \hookrightarrow B(\square)^{\otimes |\lambda|}$

Type A_r : $\boxed{1} \xrightarrow{1} \boxed{2} \xrightarrow{2} \dots \xrightarrow{r-1} \boxed{r} \xrightarrow{r} \boxed{r+1}$

Example

Type A_3

$$\begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline 4 \\ \hline \end{array} \mapsto \boxed{4} \otimes \boxed{3} \otimes \boxed{1}$$

- strictly increasing in columns

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Type C_3 $\begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \bar{3} \\ \hline \end{array} \mapsto \boxed{\bar{3}} \otimes \boxed{3} \otimes \boxed{1}$

- alphabet $[\bar{r}] := \{1 < 2 < \dots < r < \bar{r} < \overline{r-1} < \dots < \bar{1}\}$
- strictly increasing in columns
- for column $b = b(k) \dots b(1)$ there is no pair (z, \bar{z}) s.t.:

$$z = b(p), \quad \bar{z} = b(q), \quad q - p \leq k - z.$$

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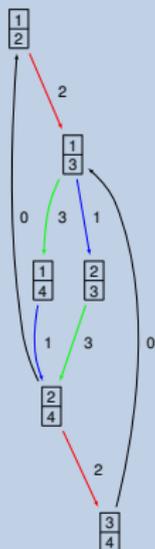
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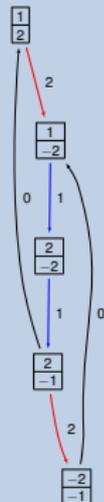
Column KR crystals for types $A_n^{(1)}$ and $C_n^{(1)}$

Example

$B^{2,1}$ of type $A_3^{(1)}$



$B^{2,1}$ of type $C_2^{(1)}$



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Crystals

Energy function

Charge

Arbitrary type

Energy function

$B := B_\mu = B^{\mu'_1, 1} \otimes B^{\mu'_2, 1} \otimes \dots$, connected by f_0 arrows.

The **energy** $D : B \rightarrow \mathbb{Z}$ originates from exactly solvable lattice models (computed via **local energies** and the **combinatorial R -matrix**).

Alternative construction (S., Tingley) as **affine** grading on B :

- constant on classical components (f_0 arrows removed)
- increases by 1 along f_0 arrows which are not at the end of a 0-string (**Demazure arrows**)

Remark

In most cases, B is still connected upon removal of non-Demazure f_0 arrows.

$\Rightarrow D$ is well-defined up to constant.

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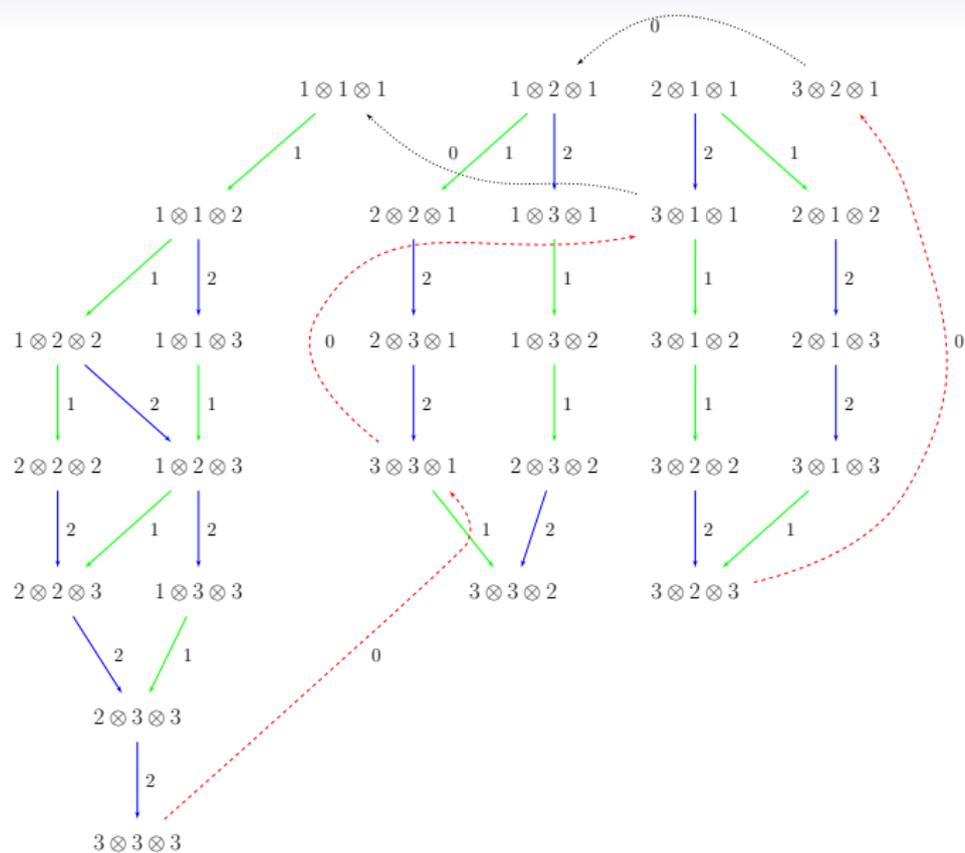
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Arbitrary type

Charge type A

Charge à la [Lascoux](#) and [Schützenberger](#):
 w word of partition content μ

Example

$$\mu = (3, 3, 3, 1)$$

1132214323

$$\text{charge}(1132214323) = 1 + 2 + 3 = 6$$

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Charge on KN tableaux - type A

$$B_\mu := \bigotimes_{i=1}^{\mu_1} B^{\mu'_i, 1}$$

circular order \prec_i : $i \prec_i i+1 \prec_i \dots \prec_i n \prec_i 1 \prec_i \dots \prec_i i-1$
 construct reordered c from $b \in B_\mu$

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$$cw(b) = \begin{pmatrix} 6 & 5 & 4 & 4 & 3 & 3 & 2 & 2 & 2 & 1 \\ 1 & 1 & 3 & 2 & 2 & 1 & 4 & 3 & 2 & 3 \end{pmatrix}$$

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$$\sum_{\gamma \in \text{Des}(c)} \text{arm}(\gamma) = \text{charge}(cw_2(b))$$

Remark

A similar construction works in type C.

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Relation between charge and energy

Theorem (Lenart, S. 2011)

$B = B^{r_N,1} \otimes \dots \otimes B^{r_1,1}$ of type $A_n^{(1)}$ or type $C_n^{(1)}$

Then for $b \in B$

$$D(b) = \text{charge}(b)$$

Idea of proof: Verify that charge satisfies the recursive relations of the energy function.

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Generalizing the charge to arbitrary root systems

Key concept: quantum Bruhat graph (QBG).

In type A_{n-1} , it is the graph on S_n with directed edges

$$w \longrightarrow wt_{ij},$$

where

$$\begin{aligned} \ell(wt_{ij}) &= \ell(w) + 1 \quad (\text{Bruhat graph}), \quad \text{or} \\ \ell(wt_{ij}) &= \ell(w) - \ell(t_{ij}) = \ell(w) - 2(j - i) + 1. \end{aligned}$$

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Generalizing the charge to arbitrary root systems

Key concept: quantum Bruhat graph (QBG).

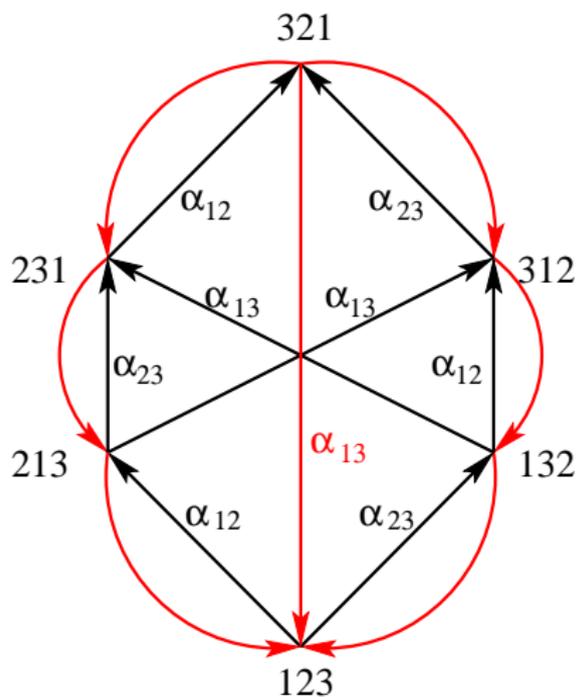
In type A_{n-1} , it is the graph on S_n with directed edges

$$w \longrightarrow wt_{ij},$$

where

$$\begin{aligned} \ell(wt_{ij}) &= \ell(w) + 1 \quad (\text{Bruhat graph}), \quad \text{or} \\ \ell(wt_{ij}) &= \ell(w) - \ell(t_{ij}) = \ell(w) - 2(j - i) + 1. \end{aligned}$$

Quantum Bruhat graph for S_3 :



The key ingredient

Fact. Fix two column strict fillings (in type A)

$$\begin{array}{|c|} \hline a_1 \\ \hline a_2 \\ \hline \vdots \\ \hline a_k \\ \hline \end{array} \quad \text{and} \quad \begin{array}{|c|} \hline b_1 \\ \hline b_2 \\ \hline \vdots \\ \hline b_k \\ \hline \end{array},$$

where the second one is reordered according to the first.

There is a unique path in the quantum Bruhat graph of the following form:

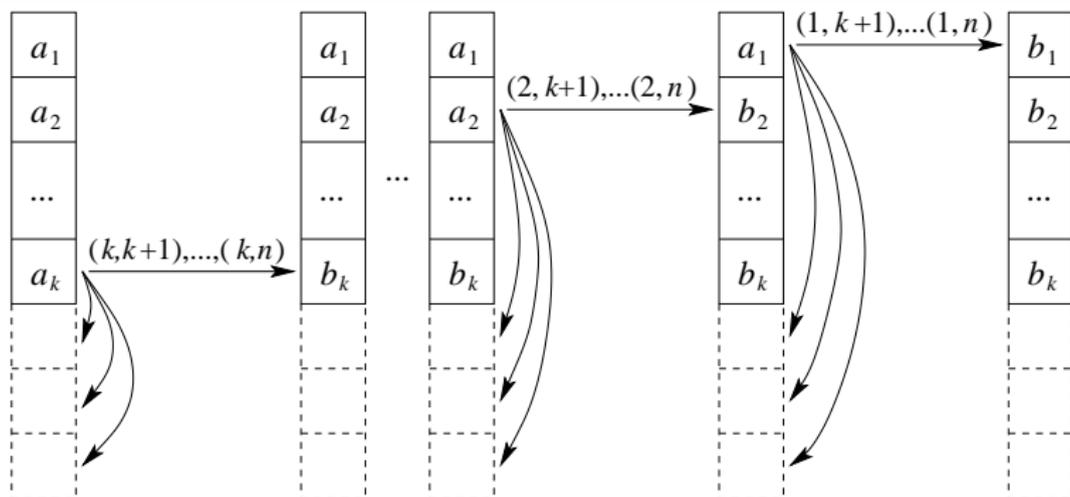
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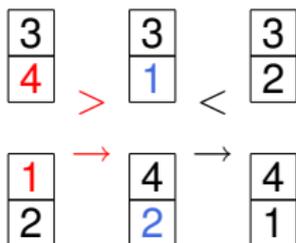
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Fillings as chains of permutations

$$b = \begin{array}{|c|c|c|c|} \hline 3 & 2 & 1 & 2 \\ \hline 4 & 3 & & \\ \hline \end{array} \mapsto c = \begin{array}{|c|c|c|c|} \hline 3 & 3 & 1 & 2 \\ \hline 4 & 2 & & \\ \hline \end{array} \mapsto \Pi = (\pi_1, \pi_2, \dots).$$



$((2, 3), (2, 4), (1, 3), (1, 4))$ |

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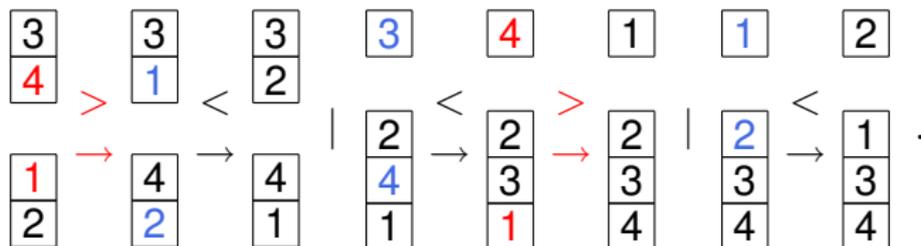
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$$\begin{array}{|c|} \hline 3 \\ \hline 4 \\ \hline \end{array} > \begin{array}{|c|} \hline 3 \\ \hline 1 \\ \hline \end{array} < \begin{array}{|c|} \hline 3 \\ \hline 2 \\ \hline \end{array} \\ \hline \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array} \xrightarrow{\text{red}} \begin{array}{|c|} \hline 4 \\ \hline 2 \\ \hline \end{array} \xrightarrow{\text{blue}} \begin{array}{|c|} \hline 4 \\ \hline 1 \\ \hline \end{array} \quad |$$

$$((2, 3), (2, 4), (1, 3), (1, 4) \mid$$

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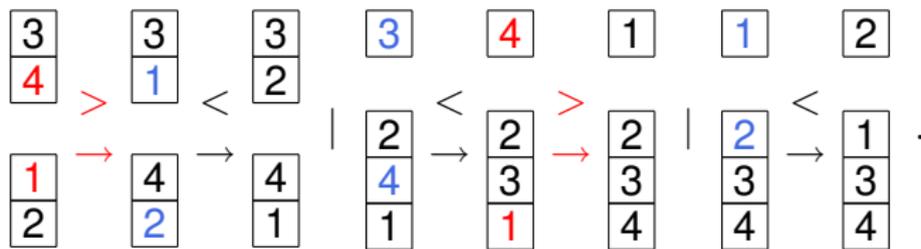
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$$l_r \quad ((2, 3), (2, 4), (1, 3), (1, 4) \mid (1, 2), (1, 3), (1, 4) \mid (1, 2), (1, 3), (1, 4))$$

1 1 3 3 2 2 2 1 1 1

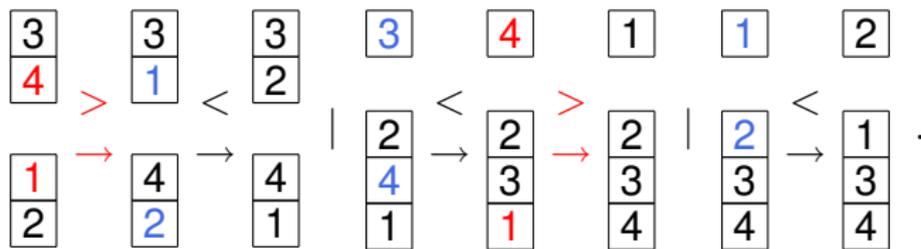


$$l_r = \text{arm}(\text{descent})$$

$$\text{charge}(b) = \sum_{\gamma \in \text{Des}(c)} \text{arm}(\gamma) = \sum_{\pi_r > \pi_{r+1}} l_r =: \text{level}(\Pi).$$

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3	3	1	2
4	2*	-	-
	*	-	-

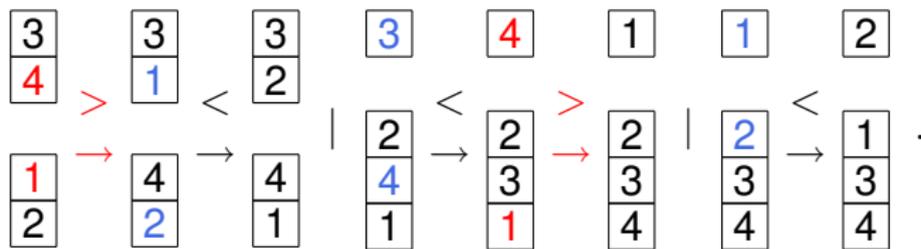
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Construction of level statistic

Step 1. Fix a partition μ .

Step 2. Associate with μ a sequence (μ -chain) Γ of pairs (i_r, j_r) (i.e., roots in type A) – several choices possible, but not explained.

Example. For $\mu = (4, 2, 0)$, we considered

$$\Gamma = ((2, 3), (2, 4), (1, 3), (1, 4) | (1, 2), (1, 3), (1, 4) | (1, 2), (1, 3), (1, 4)).$$

Step 3. Define $l_r = \#\{s \geq r : (i_s, j_s) = (i_r, j_r)\}$.

Step 4. Define **admissible subsets**:

$$\mathcal{A}(\Gamma) = \mathcal{A}(\mu) = \#\{\text{subsets } \Pi \text{ of } \Gamma \text{ giving rise to paths in the QBG}\}.$$

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Remarks.

1. The above construction works for any finite root system, as all the ingredients apply to the general case.
2. The level statistic originates in the **Ram-Yip formula** for **Macdonald polynomials** of arbitrary type:

$$(*) \quad P_{\mu}(x; q, 0) = \sum_{\Pi \in \mathcal{A}(\mu)} q^{\text{level}(\Pi)} x^{\text{weight}(\Pi)}.$$

In fact, we can rewrite (*) via the bijection between $\mathcal{A}(\mu)$ and fillings explained before (which also works in type C).

Theorem (L.)

In types A and C, we have

$$P_{\mu}(x; q, 0) = \sum_{b \in B^{\mu'_1, 1} \otimes B^{\mu'_2, 1} \otimes \dots} q^{\text{charge}(b)} x^{\text{weight}(b)}.$$

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Construction. (L. and Lubovsky) On $\mathcal{A}(\mu)$ was defined the structure of an affine crystal (purely combinatorially) – the *quantum alcove model*.

Conjecture. (L. and Lubovsky)

1. There is a bijection between $\mathcal{A}(\mu)$ in type X_n and the KR crystal $B_\mu := B^{\mu'_1,1} \otimes B^{\mu'_2,1} \otimes \dots$ of type $X_n^{(1)}$ under which the arrows of $\mathcal{A}(\mu)$ correspond to arrows of B_μ .
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- The KR crystal and its energy function are realized in terms of quantum Lakshmibai-Seshadri (LS) paths.
- For μ regular (in type A : partitions with distinct parts), the quantum LS paths are in bijection with $\mathcal{A}(\Gamma)$ for a special μ -chain Γ . The conjecture is verified in this case.
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